RESEARCH OF VEHICLES DIRECTIONAL STABILITY

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Abstract: In this work for the research of vehicles directional stability was used three different mechanical-mathematical models. The behavior of car stability has been simulated with MATLAB software.

Keywords: DIRECTIONAL STABILITY, MECHANICAL-MATHEMATICAL MODEL, LATERAL ELASTICITY, SLIP ANGLE

1. Introduction

Stability of a vehicle concerns itself with the tendency of a vehicle to return to its original direction when disturbed (rotated) away from that original direction.

The goal of this work is research of vehicles directional stability. The situation is described by three different mechanical-mathematical models around to the axes O_v and O_z .

During the research we make the following considerations:

- The car and his models are symmetrical to the axe O_x
- The dynamical process (displacement around to the axes O_x) isn't exanimated
- The main factor of cars stability is the wheels slip angle.
- Working in the zone of pure rolling motion (the slip angle for the wheels was zero). Figure 1 shows the general relationship between the lateral force a tire F_y and the slip angle of the tire δ_y .



Fig. 1 Relationship between the lateral force a tire F_y and the slip angle of the tire δ_y

Mechanical-mathematical model examine only the lateral elasticity

The behavior of the car is described using mechanicalmathematical model. Scheme of the model is shown to Figure 2 and the suspended masses include the masses of the elements of the car, passengers and load.. If we consider that the slip angle for the wheels was zero we could examine the car as a mass witch point 1, 2, 3 and 4 around to the axe O_y are bending.

The motion of system is exanimated as function to the displacements are around the axes O_y and the angular displacement around the axes O_z .

In the center of gravity is fixed local coordinate system attached $O_0x_0y_0z_0$. All displacements of local coordinate systems are given to the absolute coordinate system $O_Ax_Ay_Az_A$. In the equilibrium position the axis of all coordinate systems are parallel.



Fig. 2 Mechanical-mathematical model, examined the zone of directional stability

To find the lows motion to the absolute coordinate system $O_A x_A y_A z_A$ is necessary to define the transition matrices of each local coordinate systems to the absolute.

For generalized coordinate systems are assumed:

- y₀ linear displacement of the local coordinate system O₀x₀y₀z₀ to absolute O_Ax_Ay_Az_A around axis O_y;
- ψ₀ angular displacement of the coordinate system
 O₀x₀y₀z₀ to absolute O_Ax_Ay_Az_A around axis O_z;

Matrix of transition from $O_0 x_0 y_0 z_0$ to $O_A x_A y_A z_A$ to the system is:

$$T_{o}^{A} = \begin{bmatrix} \cos\psi_{o} & -\sin\psi_{o} & 0 & 0\\ \sin\psi_{o} & \cos\psi_{o} & 0 & y_{0}\\ 0 & 0 & 1 & 0\\ 0 & 0 & 0 & 1 \end{bmatrix}$$

 x_0 and z_0 are zero (we consider only linear displacement around axis O_y , i.e. only laterally).

$$T^A$$

The matrix of transition ${}^{I_{o}}$ can be simplified with the assuming that ψ_0 is small witch leads to the linearization of trigonometrically functions ($\sin\psi_0 \approx \psi_0$ and $\cos\psi_0 \approx 1$):

$$T_o^A = \begin{bmatrix} 1 & -\psi_o & 0 & 0 \\ \psi_o & 1 & 0 & y_0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

Radius vector of the point gravity: $\rho_{mo} = [0, y_0, 0, 1]^T$

Radius vector of the point 1: $\rho_1^0 = [a_1, -b, 0, 1]^T$

$$\rho_{1}^{A} = T_{0}^{A} \cdot \rho_{1}^{0} \rightarrow \begin{pmatrix} a_{1} + b\psi_{0} \\ y_{0} + a_{1}\psi_{0} - b \\ 0 \\ 1 \end{pmatrix}$$

Radius vector of the point 2: $\rho_2^0 = [a_2, b, 0, 1]^T$

$$\rho_{2}^{A} = T_{0}^{A} \cdot \rho_{2}^{0} \to \begin{pmatrix} a_{1} - b\psi_{0} \\ y_{0} + a_{1}\psi_{0} + b \\ 0 \\ 1 \end{pmatrix}$$

Radius vector of the point 3: $\rho_3^0 = [-a_2, -b, 0, 1]^T$

$$\rho_{3}^{A} = T_{0}^{A} \cdot \rho_{3}^{0} \rightarrow \begin{pmatrix} b \psi_{0} - a_{2} \\ y_{0} - a_{2} \psi_{0} - b \\ 0 \\ 1 \end{pmatrix}$$

Radius vector of the point 4: $\rho_4^0 = [-a_2, b, 0, 1]^T$

$$\rho_{4}^{A} = T_{0}^{A} \cdot \rho_{4}^{0} \rightarrow \begin{pmatrix} -a_{2} - b\psi_{0} \\ y_{0} - a_{2}\psi_{0} + b \\ 0 \\ 1 \end{pmatrix}$$

The components of the angular velocity of the system are set in advance:

$$\omega_{OX}^{A} = 0$$

$$\omega_{OY}^{A} = 0$$

$$\omega_{OZ}^{A} = \dot{\psi}$$

erms is

$$T = \frac{1}{2}m_{0}\dot{y}_{0}^{2} + \frac{1}{2}J_{oz}\dot{\psi}_{0}^{2}$$

Kinetic energie of the systems is

Potential energie of the systems is

$$\Pi = \frac{1}{2}c_{y11}(y_0 + a_1\psi_0)^2 + c_{y12}(y_0 + a_1\psi_0)^2 + c_{y21}(y_0 - a_2\psi_0)^2 + c_{y22}(y_0 - a_2\psi_0)^2$$

Appling Lagrange's equation of 2nd kind

$$\frac{d}{dt}\left(\frac{\partial T}{\partial \dot{q}}\right) - \left(\frac{\partial T}{\partial q}\right) = -\left(\frac{\partial II}{\partial q}\right) - \left(\frac{\partial R}{\partial \dot{q}}\right)$$

The differentials equations of the system are

$$\begin{split} m_{0}\ddot{y}_{o} + & \left(c_{y11} + c_{y12} + c_{y21} + c_{y22}\right)y_{o} + \left(\left(c_{y11} + c_{y12}\right)a_{1} - \left(c_{y21} + c_{y22}\right)a_{2}\right)\psi_{0} = F_{y} \\ J_{oz}\ddot{\psi}_{o} + & \left(\left(c_{y11} + c_{y12}\right)a_{1}^{2} + \left(c_{y21} + c_{y22}\right)a_{2}^{2}\right)\psi_{o} + & \left(\left(c_{y11} + c_{y12}\right)a_{1} - \left(c_{y21} + c_{y22}\right)a_{2}\right)y_{0} = M_{y} \end{split}$$

For the differentials equations describing the system from figure 2 are valid:

$$[M]\ddot{q} + [B]\dot{q} + [C]q = [F]_{, \text{ където}}$$

[M] is the matrix of inertia witch is symmetrical to the main diagonal, with dimension 2x2 and has the following form:

$$\begin{bmatrix} M \end{bmatrix} = \begin{pmatrix} m_0 & 0 \\ 0 & J_{oz} \end{pmatrix}$$

[C] is the matrix of elasticity which is also symmetric to the main diagonal and has dimension 2x2:

$$[C] = \begin{pmatrix} (c_{y11} + c_{y12} + c_{y21} + c_{y22}) & ((c_{y11} + c_{y12})a_1 - (c_{y21} + c_{y22})a_2) \\ ((c_{y11} + c_{y12})a_1 - (c_{y21} + c_{y22})a_2) & ((c_{y11} + c_{y12})a_1^2 + (c_{y21} + c_{y22})a_2^2) \end{pmatrix}$$

[B] is the matrix of dissipative forces, showing the influence of damper. Also symmetric with dimension 2x2, in our case [B] = 0.

The generalized coordinates and their derivatives are:

$$q\} = \begin{bmatrix} y_0 \\ \psi_0 \end{bmatrix}; \{\dot{q}\} = \begin{bmatrix} \dot{y}_0 \\ \dot{\psi}_0 \end{bmatrix}; \{q\} = \begin{bmatrix} \ddot{y}_0 \\ \ddot{\psi}_0 \end{bmatrix};$$

To obtain natural frequencies of the system, the equations are represented in Cauchys' normal form: y+Ly = 0, where L has the following form:

$$L = \begin{bmatrix} M^{-1}B & M^{-1}C \\ I & 0 \end{bmatrix}$$

The output parameters of the system – displacement, velocity and acceleration are obtained from the equations: y+Ly = Y, where Y is:

$$Y = \begin{bmatrix} M^{-1} & F(t) \\ 0 \end{bmatrix}$$

All solutions to the definite time interval are obtained after integration of the system using the method of Runge - Kutta.

Plane model examine only the slip angle

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Figure 3 shows the schema of the model exanimating the car stability with reporting the wheel slip angel. The model is known as Rocards model.



Fig. 3 Plane model describing only the slip angle

The initial conditions were assumed for modeling:

- The point 1, 2, 3 and 4 to the fig.3 are the projection of the wheels;
- Considering that the wheels are undeformables and can ignore their lateral elasticity around to the axis O_y.
- The car is shown at the angle ψ to the axis O_x.

From fig. 4 the slip angles are δ_1 , δ_2 , δ_3 , δ_4 can be determinate by coordinate y, so we have:



Fig. 4 Slip angle

$$tg(\psi) = \frac{dy}{dx} \to \psi = \frac{dy}{dx} \to dy = \psi.dx$$
$$tg(\psi + \delta) = \frac{dy_1}{dx} \to \delta = \frac{dy_1}{dx} - \psi \to \delta = \frac{dy_1 - dy}{dx} = \frac{dy_1 - \psi.dx}{v.dt}$$

If the vehicle moved with a pure rolling motion (the slip angle for the wheels were all zero) then the displacements of the point of contact of the wheels with the ground would be the same as the displacement of the center of gravity C in any interval δt : $dx=dx_1=dx_2=dx_3=dx_4=V.dt$, where V is the speed of the vehicle in direction O_x .

The equations of motion are:

$$m\ddot{y} + k_1(\delta_1 + \delta_2) + k_2(\delta_3 + \delta_4) = F_y$$
$$J_{OZ}\ddot{\psi} + k_1a_1(\delta_1 + \delta_2) - k_2a_2(\delta_3 + \delta_4) = M_y$$

Substituting for δ_1 , δ_2 , δ_3 , δ_4 we obtain:

$$\begin{split} m\ddot{y} &+ \frac{2}{v}(k_1 + k_2)\dot{y} + \frac{2}{v}(k_1a_1 - k_2a_2)\dot{\psi} - 2(k_1 + k_2)\psi = F_y \\ J_{oz}\ddot{\psi} &+ \frac{2}{v}(k_1a_1^2 + k_2a_2^2)\dot{\psi} - 2(k_1a_1 - k_2a_2)\psi + \frac{2}{v}(k_1a_1 - k_2a_2)\dot{y} = M_y \end{split}$$

Taking F_y , M_{ψ} to be zero, $I=m.\rho^2$ and after substituting these values to the equations of motion, we are searching the solutions in the form $y=Ae^{pt}$ and $\psi=Be^{pt}$. To have a solution the equations system is necessary the matrix from the coefficient in front of the A and B to be equal to zero. So, we have:

$$\begin{vmatrix} p^{2} + 2\frac{(k_{1} + k_{2})}{mv}p & 2\frac{(k_{1}a_{1} - k_{2}a_{2})}{mv}p - 2\frac{(k_{1} + k_{2})}{mv}\\ 2\frac{(k_{1}a_{1} - k_{2}a_{2})}{mv}p & p^{2} + 2\frac{(k_{1}a_{1}^{2} + k_{2}a_{2}^{2})}{m\rho^{2}v}p - \frac{(k_{1}a_{1} - k_{2}a_{2})}{m\rho^{2}} \end{vmatrix} = 0$$

After matrix development is obtained:

$$p^{4} + \frac{2}{m\nu} \left[k_{1} \left(1 + \frac{a_{1}^{2}}{\rho^{2}} \right) + k_{2} \left(1 + \frac{a_{2}^{2}}{\rho^{2}} \right) \right] p^{3} + \left[\frac{4 \left[\left(k_{1} + k_{2} \right) \left(k_{1}a_{1}^{2} + k_{2}a_{2}^{2} \right) - \left(k_{1}a_{1} - k_{2}a_{2} \right)^{2} \right] - \frac{2 \left(k_{1}a_{1} - k_{2}a_{2} \right)}{m\rho^{2}} \right] p^{2} = 0$$

This may be divides by p^2 :

$$p^{2} + \frac{2}{mv} \left[k_{1} \left(1 + \frac{a_{1}^{2}}{\rho^{2}} \right) + k_{2} \left(1 + \frac{a_{2}^{2}}{\rho^{2}} \right) \right] p + \left[-\frac{2(k_{1}a_{1} - k_{2}a_{2})}{m\rho^{2}} + \frac{4k_{1}k_{2}(a_{1} + a_{2})^{2}}{m^{2}\rho^{2}v^{2}} \right] = 0$$

The analysis of stability is made with the criteria of Raus-Kurvits by creating a square matrix from the characteristic equation $A(s) = a_0 s^n + a_1 s^{n-1} + ... + a_{n-1} s + a_n = 0$.

For being systems stability is necessary $a_0 > 0$ and $\Delta_1, \Delta_2, ..., \Delta_n$ to be always positive. Where $\Delta_1, \Delta_2, ..., \Delta_n$ are:

$$\Delta_1 = \|a_1\|, \Delta_2 = \|a_1 - a_3\|, \Delta_3 = \|a_1 - a_3 - a_5\|, \dots, \Delta_n = \|\Gamma\|$$

In our case $a_0 > 0$, so to have a cars stability motion is necessary the moment of intertie of front wheels to be smaller than that of rear wheels.

For the vehicles speed is obtained:

$$v^{2} \geq \frac{\left\lfloor \frac{4k_{1}k_{2}(a_{1}+a_{2})^{2}}{m^{2}\rho_{2}} \right\rfloor}{\frac{2(k_{1}a_{1}-k_{2}a_{2})}{m\rho^{2}}} = \frac{2k_{1}k_{2}(a_{1}+a_{2})^{2}}{m[k_{1}a_{1}-k_{2}a_{2}]}$$

Plane model examine lateral elasticity and slip angle

The initial conditions for this case of modeling were the same as those from fig.3. But this time the lateral elasticity of the wheel is added (fig.5).



Fig. 5 Plane model about lateral elasticity and slip angle

Appling the models initial condition for the differentials equations describing the system are valid:

$$\begin{split} m\ddot{y} + k_1(\delta_1 + \delta_2) + k_2(\delta_3 + \delta_4) + 2(c_1 + c_2)y + 2(c_1a - bc_2)\psi &= F_y \\ J_{oZ}\ddot{\psi} + k_1a_1(\delta_1 + \delta_2) - k_2a_2(\delta_3 + \delta_4) + 2(c_1a^2 + c_2b^2)\psi + 2(c_1a - bc_2)y &= M_y \end{split}$$

Substituting for δ_1 , δ_2 , δ_3 , δ_4 we obtain:

$$m\ddot{y} + \frac{2}{\nu}(k_1 + k_2)\dot{y} + \frac{2}{\nu}(k_1a_1 - k_2a_2)\dot{y} - 2(k_1 + k_2)\psi + 2(c_1 + c_2)y + 2(c_1a - bc_2)\psi = F_y$$

$$J_{OZ}\ddot{\psi} + \frac{2}{\nu}(k_1a_1^2 + k_2a_2^2)\dot{\psi} - 2(k_1a_1 - k_2a_2)\psi + \frac{2}{\nu}(k_1a_1 - k_2a_2)\dot{y} + 2(c_1a^2 + c_2b^2)\psi + 2(c_1a - bc_2)y = M_y$$

1. Computer Simulation

All vehicles models permit by variable of the input information to explore the influence of some construction parameters into the directional stability.

Vehicle characteristics and parameters' numerical values were taken directly from the literature.

To solve the equations was made a program with MATLAB software. It simulates the behavior of car with different speed, with different position of center of gravity and variable characteristic of wheels. The investigations were made in two parts – without and with disturbing force. This force works only for a few seconds.

m	Vehicle mass	1 500	kg
а	Distance from CG to front axel	1	m
b	Distance from CG to rear axel	1,5	m
I _{z0}	Moment of inertia to axe Oz	2 500	Kgm2
cy	Lateral stiffness of each tire to axe Oy	57 000	N/m

2. Results

Simulation results of the first model - displacements around the axes O_y and the angular displacement around the axes O_z when the center of gravity changes the position and for different value of lateral elasticity are given in table 1.

c _y , kN/m		Distance from CG to front axel a ₁ , m					
Front/rear wheel		1	1,2	1,25	1,3	1,5	
50/50	Displacement around Oy, m	0,008714	0,008624	0,008658	0,008632	0,008694	
	Angular displacement around O _z , grad	0,3924	0,3116	0,008931	0,2995	0,3895	
60/60	Displacement around Oy, m	0,00796	0,007885	0,007903	0,007869	0,007943	
	Angular displacement around O _z , grad	0,3586	0,2845	0,008154	0,2734	0,3556	
70/70	Displacement around Oy, m	0,007391	0,007312	0,007311	0,007286	0,007354	
	Angular displacement around O ₂ , grad	0,3321	0,2634	0,007553	0,2531	0,3288	
80/80	Displacement around Oy, m	0,006928	0,006847	0,006843	0,006841	0,006898	
	Angular displacement around O _z , grad	0,3105	0,2463	0,007066	0,2363	0,3076	
90/90	Displacement around Oy, m	0,006534	0,006458	0,006454	0,006452	0,006487	
	Angular displacement around O _z , grad	0,2921	0,2323	0,006661	0,2232	0,2902	
50/60	Displacement around Oy, m	0,00831	0,008249	0,008232	0,008243	0,008256	
	Angular displacement around O _z , grad	0,37	0,3661	0,3519	0,3111	0,3537	
60/70	Displacement around Oy, m	0,007809	0,007592	0,007594	0,007583	0,007582	
	Angular displacement around O ₂ , grad	0,3459	0,3331	0,3181	0,2553	0,3329	
70/80	Displacement around Oy, m	0,007264	0,007076	0,007071	0,007067	0,00706	
	Angular displacement around O ₂ , grad	0,3221	0,3101	0,2896	0,202	0,3084	
80/90	Displacement around Oy, m	0,006807	0,06661	0,006646	0,006642	0,006637	
	Angular displacement around O _z , grad	0,3009	0,2881	0,2655	0,1515	0,2936	
90/100	Displacement around Oy, m	0,006364	0,006292	0,006287	0,006284	0,006312	
	Angular displacement around O _z , grad	0,2859	0,27	0,2443	0,1057	0,2784	
60/50	Displacement around Oy, m	0,008287	0,008219	0,008218	0,008231	0,008253	
	Angular displacement around O _z , grad	0,3611	0,3021	0,3479	0,3599	0,3657	
70/60	Displacement around Oy, m	0,007619	0,007577	0,007578	0,007573	0,007756	
	Angular displacement around O _z , grad	0,3365	0,2442	0,3127	0,3307	0,345	
80/70	Displacement around Oy, m	0,00708	0,007063	0,007061	0,00706	0,007225	
	Angular displacement around O _z , grad	0,3154	0,19	0,2832	0,3052	0,322	
90/80	Displacement around Oy, m	0,006658	0,006639	0,006637	0,006637	0,006774	
	Angular displacement around O _z , grad	0,2976	0,1393	0,2576	0,2851	0,3008	
100/90	Displacement around Oy, m	0,006332	0,006282	0,00628	0,00628	0,006332	
	Angular displacement around O _z , grad	0,2815	0,09346	0,2376	0,2675	0,2844	

Simulations results of the second model shown in fig. 6...9 demonstrate the displacements are around the axes O_y and the angular displacement around the axes O_z as function of coordinate of center of gravity, the steering force characteristics for the front k_1 and rear tires k_2 and vehicles speed.



fig. 6 Displacements are around the axes O_y and the Angular displacement around the axes O_z Distance from CG to front axel $a_1 = 1.2$,



fig. 7 Displacements are around the axes O_y and the Angular displacement around the axes O_z Distance from CG to front axel $a_1 = 1.2$, $V_{cr} = 0 \text{ km} / h$



fig. 8 Displacements are around the axes O_y and the Angular displacement around the axes $O_z\,$ Distance from CG to front axel a_1 = 1.3, V_{cr} = 294 km / h



fig. 9 Displacements are around the axes O_y and the Angular displacement around the axes O_z Distance from CG to front axel a_1 = 1.3, V_{cr} = 294 km / h

Simulations results of the third model shown in fig. 10...17 demonstrate the displacements are around the axes O_y and the angular displacement around the axes O_z as function of coordinate of center of gravity, the steering force characteristics for the front k_1 and rear tires k_2 and vehicles speed.



fig. 10 Displacements are around the axes O_y and the Angular displacement around the axes O_z Distance from CG to front axel a_1 = 1.2, $V_{\rm cr}$ = 0 km / h



fig. 11 Displacements are around the axes O_y and the Angular displacement around the axes O_z Distance from CG to front axel a_1 = 1.2, $V_{\rm cr}$ = 0 km / h



fig. 12 Displacements are around the axes O_y and the Angular displacement around the axes $O_z\,$ Distance from CG to front axel a_1 = 1.2, $V_{\rm cr}$ = 0 km / h



fig. 13 Displacements are around the axes O_y and the Angular displacement around the axes O_z Distance from CG to front axel a_1 = 1.2, $V_{\rm cr}$ = 0 km / h



fig. 14 Displacements are around the axes O_y and the Angular displacement around the axes O_z Distance from CG to front axel a_1 = 1.3, V_{cr} = 294 km / h



fig. 15 Displacements are around the axes O_y and the Angular displacement around the axes O_z Distance from CG to front axel $a_1 = 1.3$, $V_{\rm cr} = 294$ km / h



fig. 16 Displacements are around the axes O_y and the Angular displacement around the axes O_z Distance from CG to front axel $a_1 = 1.3$, $V_{cr} = 294$ km / h



fig. 17 Displacements are around the axes O_y and the Angular displacement around the axes O_z Distance from CG to front axel $a_1 = 1.3$, $V_{cr} = 294$ km / h

3. Conclusions

The analysis in this paper shows the influence of some construction parameters into the directional stability as coordinate of center of gravity, cars speed, the steering force characteristic.

4. Acknowledgement

This work is a part of the research project to support of PhD students № 122ПД0006-04 funded by Research Centre at Technical University – Sofia.

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